

# Optical Absorbance of $\beta$ -Bi<sub>4</sub>X<sub>4</sub>

Physics REU Symposium, University of Texas at Dallas

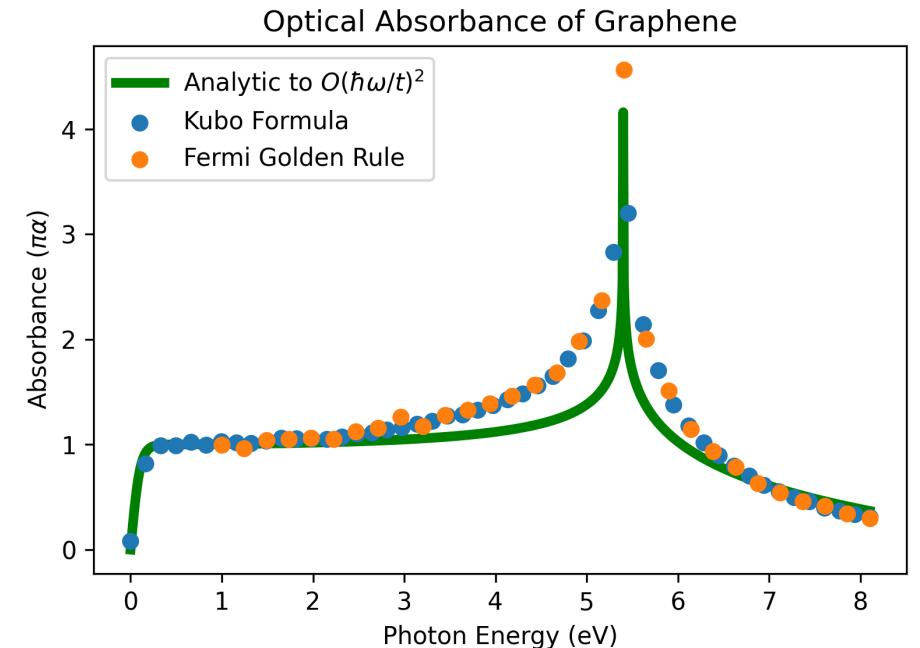
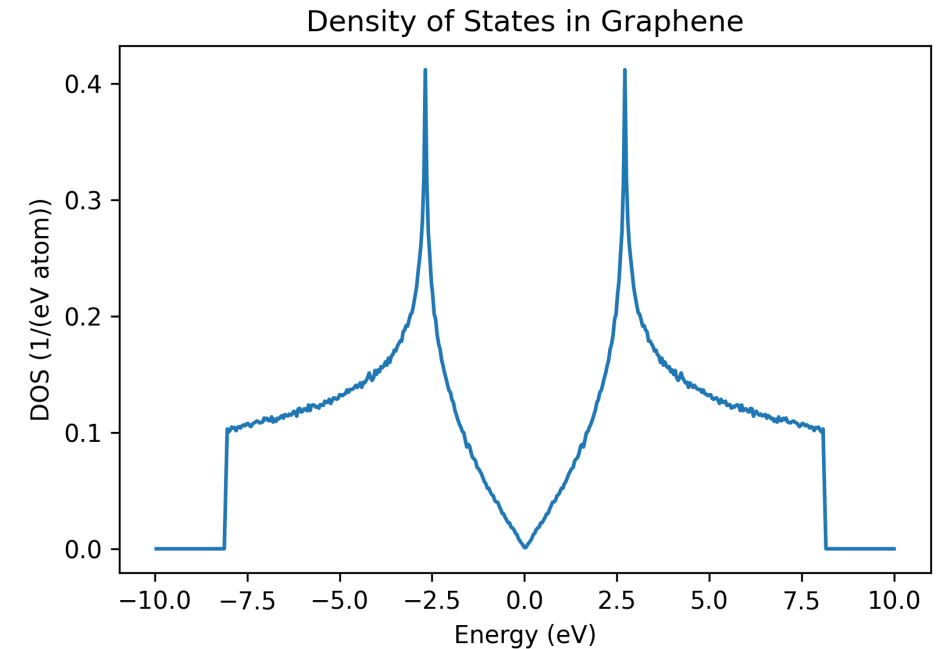
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# Optical Properties of Solids

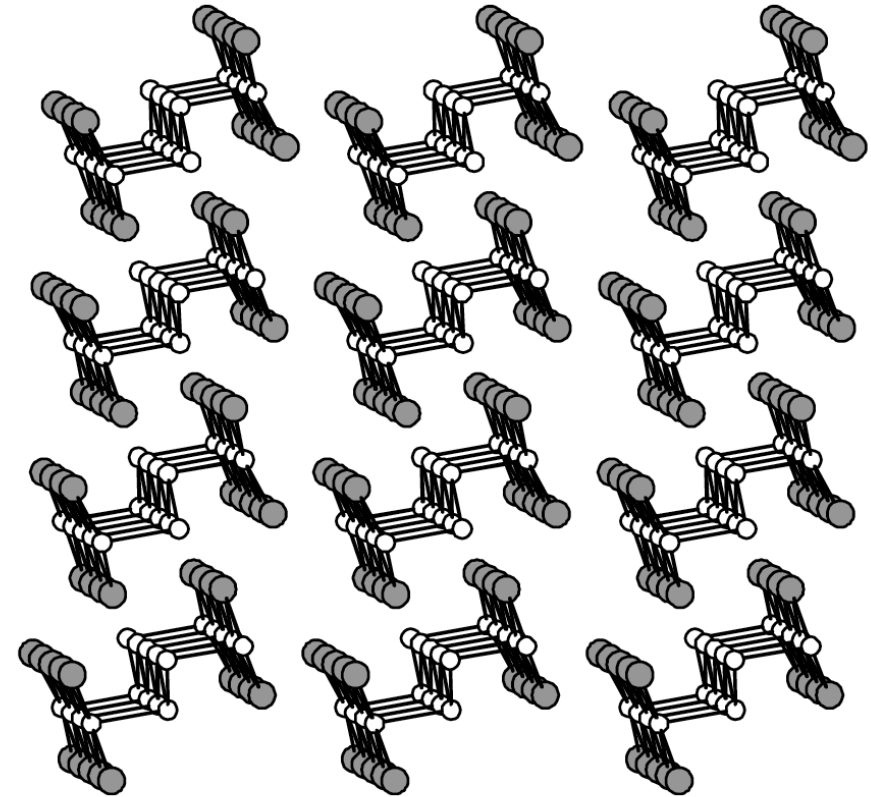
- Optical properties of solids are readily measured in experiments
- We care about optoelectronic properties because they have device applications and they help us understand materials
- Properties include absorbance, which is the amount of energy absorbed per energy incident
- Graphene's quantized absorbance<sup>1</sup>

<sup>1</sup> Nair, et al Science 320, 1308 (2008)



# Bismuth Bromide and Bismuth Iodide

- Quasi-one-dimensional van der Waals materials
  - Strands that form layers that stack
  - Protected surface states in beta phase: it is a weak topological insulator<sup>1,2</sup>
  - Protected edge states in alpha phase: it is a higher-order topological insulator<sup>3</sup>
- Two forms:
  - AA stacking (beta phase)
  - AB stacking (alpha phase)



$\beta$ -Bi<sub>4</sub>I<sub>4</sub> structure<sup>4</sup>

<sup>1</sup> Zhang et al PRL 116, 066801 (2016)

<sup>2</sup> Noguchi et al Nature 566, 518 (2019)

<sup>3</sup> Yoon, et al arXiv:2005.14710 (2020)

<sup>4</sup> Dikarev et al RCB 50(12), 2304 (2001)

# The Bulk Hamiltonian (Bloch Hamiltonian)

- In lattices (discrete translational invariance), Bloch's Theorem holds, where  $u_n(r)$  is a function with the periodicity of the lattice and  $k$  is a wavevector:

$$\psi_n(r) = \exp(ik \cdot r) u_n(r)$$

- The periodic functions  $u_n(r)$  are given by the Bulk Hamiltonian,  $\mathcal{H}(k)$ :

$$\mathcal{H}(k) u_n(r) = \epsilon_n u_n(r)$$

- We define the valence states to be:

$$v_n = \psi_n(r) \text{ such that } \epsilon_n < 0$$

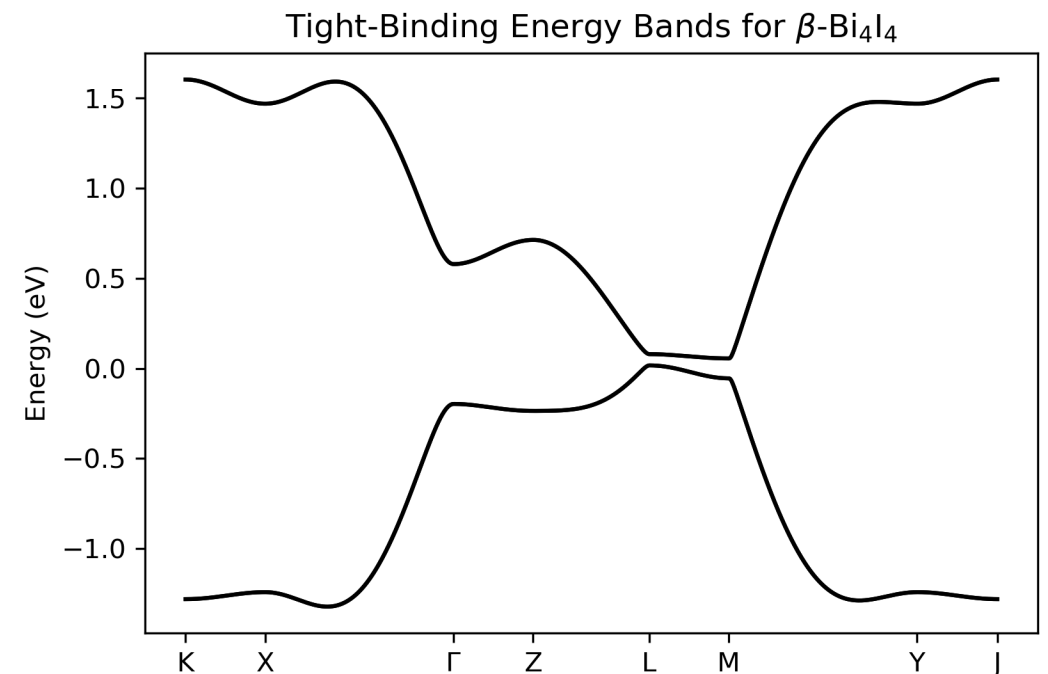
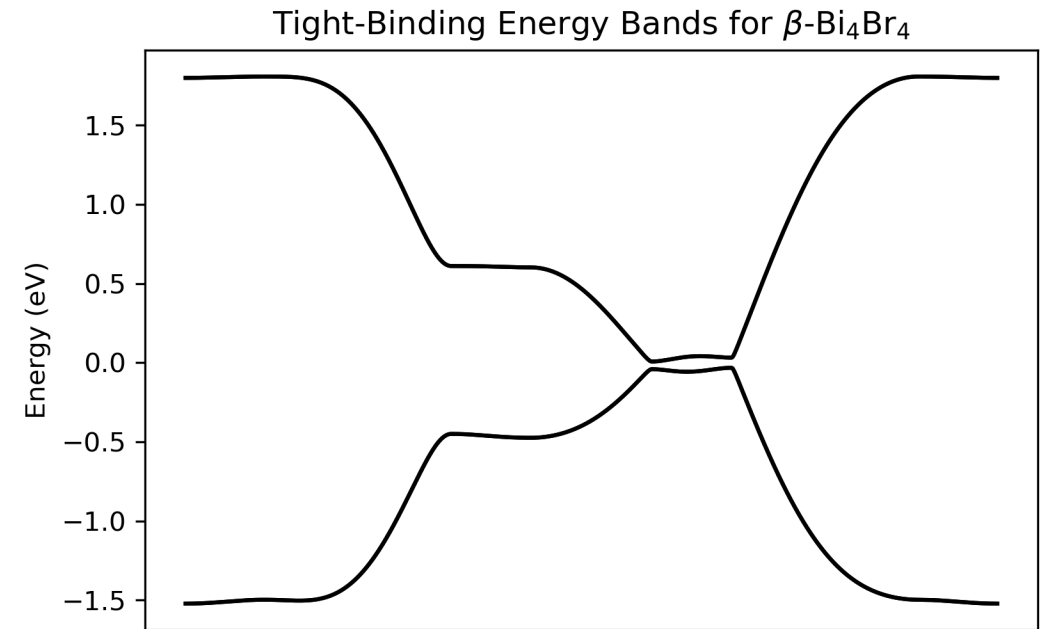
- We define the conduction states to be:

$$c_n = \psi_n(r) \text{ such that } \epsilon_n > 0$$

# Energy Bands of $\beta\text{-Bi}_4\text{X}_4$

- Energy bands: diagonalize Hamiltonian
- Tight-binding model Hamiltonian
  - Onsite energies
  - Transfer energies
  - Phase accumulation (Peierels)
- How to get a tight-binding model?
  - First principles density functional theory calculations leads to maximally localized Wannier function tight-binding models
  - Symmetry considerations and restrict to nearest neighbor transfers
  - This has been done in the literature<sup>1</sup>

<sup>1</sup> Yoon, et al arXiv:2005.14710 (2020)



# Absorbance

- Absorbance is the fraction of the incident power absorbed:

$$P = \frac{W_{\text{absorbed}}}{W_{\text{incident}}}$$

- Incident light carries energy flux (Poynting Vector, in Gaussian units):

$$W_{\text{incident}} = \frac{\omega^2 |A|^2}{4\pi c}$$

- From a conductance standpoint (Ohm's Law):

$$W_{\text{absorbed}} = \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix}^\dagger \begin{pmatrix} \sigma_{xx} & \sigma_{xy} & \sigma_{xz} \\ \sigma_{yx} & \sigma_{yy} & \sigma_{yz} \\ \sigma_{zx} & \sigma_{zy} & \sigma_{zz} \end{pmatrix} \begin{pmatrix} E_x \\ E_y \\ E_z \end{pmatrix}$$

# The Current Operator

- In electromagnetism, the current operator is  $j = \sigma E$ , or if  $E = E^\beta \hat{\beta}$ :

$$j^\alpha = \sigma_{\alpha\beta} E^\beta$$

- In the single-particle framework, the current operator is, for direction  $\alpha = x, y$ , or  $z$ :

$$j^\alpha = e v^\alpha$$

- Now, we note that for the free particle Hamiltonian:

$$\frac{\partial H_0}{\partial p} = \frac{\partial p^2/2m}{\partial p} = \frac{p}{m} = v$$

- This motivates the gradient approximation:

$$v^\alpha = \frac{\partial \mathcal{H}(k)}{\partial p^\alpha}$$

- So the current operator is, with  $p^\alpha = \hbar k^\alpha$ :

$$j^\alpha = \frac{e}{\hbar} \frac{\partial \mathcal{H}(k)}{\partial k^\alpha}$$

# Linear Response: the Kubo Formalism

- For the conductivity, integrate over the Brillouin Zone ( $k$ -space):

$$\sigma_{\alpha\beta}(\hbar\omega, \eta) = i \frac{e^2}{\hbar} \sum_{c,v} \int_{\text{BZ}} \frac{d^{\text{dim}}k}{(2\pi)^{\text{dim}}} \frac{f(\epsilon_v) - f(\epsilon_c)}{\epsilon_c - \epsilon_v} \frac{\hbar v_{vc}^\alpha \hbar v_{cv}^\beta}{\hbar\omega - (\epsilon_c - \epsilon_v) + i\eta}$$

- In the gradient approximation:

$$\hbar v_{vc}^\alpha \hbar v_{cv}^\beta = \left\langle v(k) \left| \frac{\partial \mathcal{H}(k)}{\partial k^\alpha} \right| c(k) \right\rangle \left\langle c(k) \left| \frac{\partial \mathcal{H}(k)}{\partial k^\beta} \right| v(k) \right\rangle$$

- The Fermi-Dirac distribution function is:

$$f(\epsilon) = (\exp(\epsilon/k_B T) + 1)^{-1}$$



# How to Calculate Absorbance? (the Kubo way)

- Choose system
- Find bulk Hamiltonian,  $H$
- Find eigenenergies and eigenfunctions of  $H$
- Specify polarization  $\hat{E}$ , frequency  $\omega$ , temperature  $T$ , and broadening  $\eta$
- Calculate the conductivity tensor (with Kubo)
  - At each point in the BZ, find conductivity matrix elements and evaluate; sum over these points
- Find the energy absorbed  $W(\omega) = E^\dagger \sigma(\omega) E$
- Find the absorbance  $P(\omega) = W_{\text{abs}}/W_{\text{inc}}$

Absorbance of  $\beta$ -Bi<sub>4</sub>X<sub>4</sub>

sensitive figures removed

# Conclusions

- Results

- As expected from its anisotropic structure,  $\beta$ -Bi<sub>4</sub>X<sub>4</sub> has a correspondingly anisotropic absorbance
- Absorbance is as would be expected from the density of states
- The monolayer system has substantially *lower* absorbance than the full system
- $\beta$ -Bi<sub>4</sub>X<sub>4</sub> are real materials and expect our predictions to agree with experiments

- Future directions:

- Compare and contrast alpha and beta phases
- Address other optical properties such as the photocurrent
- Consider the optical properties of surface/edge states

Questions?