

Non-Abelian Fundamental Groups, the Second Homotopy Group and Exact Sequences

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Algebraic Topology in Physics Seminar
University of Pennsylvania
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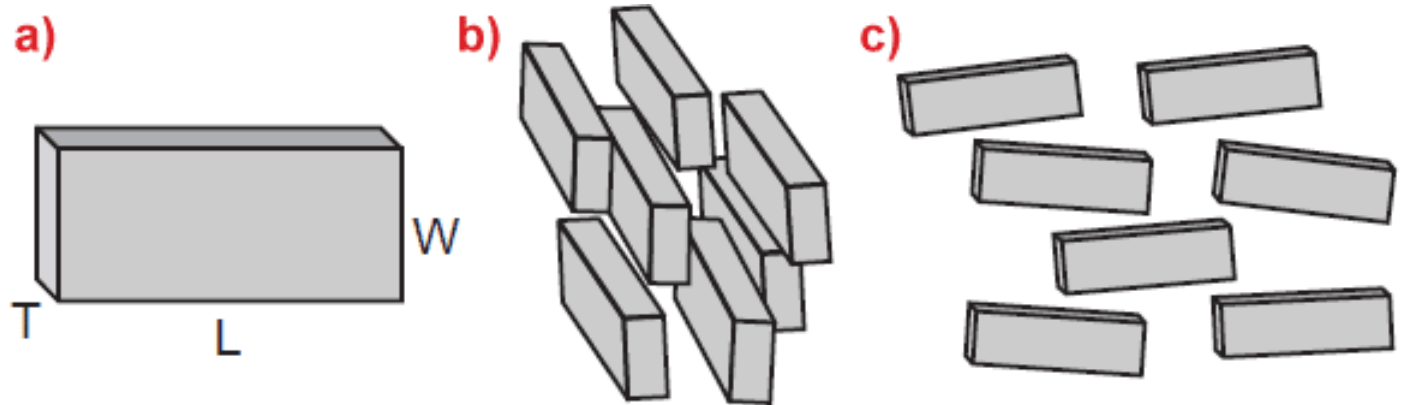
Outline

- Two (somewhat) disjoint topics
 - 1. Non-Abelian fundamental groups
 - 2a. The second homotopy group
 - 2b. Applications to quantum condensed matter physics

VI. Non-Abelian Fundamental Groups

Biaxial Nematic Liquid Crystals

- Treat as rectangular prisms
 - Identity
 - Inversion
 - π rotations (3)
 - $-\pi$ rotations (3)
- This is a group
 - The quaternion group
 - It is non-abelian



[Image Credit: ESRF](#)

Quaternions

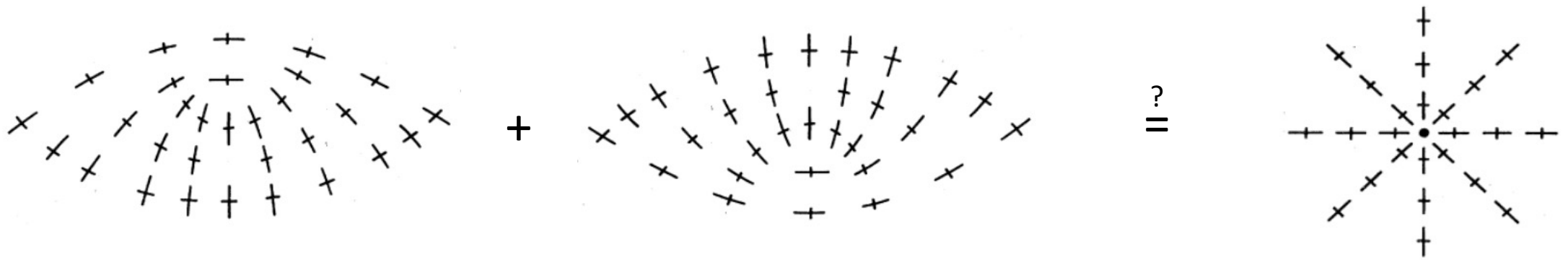
- Representation w/ Pauli matrices
 - $H = \{\pm 1, i=\pm is_x, j=\pm is_y, k=\pm is_z\}$
- Reduces to 5 conjugacy classes
 - $C_0, \bar{C}_0, C_x, C_y, C_z$
- Group multiplication table
 - Some elements are non-commuting
- Class multiplication table
 - Class multiplication is abelian so there is no notion of enclosing defects first

	1	i	-1	-i	j	k	-j	-k
1	1	i	-1	-i	j	k	-j	-k
i	i	-1	-i	1	j	-j	-k	j
-1	-1	-i	1	i	-j	-k	j	k
-i	-i	1	i	-1	-k	j	k	-j
j	j	-k	-j	k	-1	i	1	-i
k	k	j	-k	-j	-i	-1	i	1
-j	-j	k	j	-k	1	-i	-1	i
-k	-k	-j	k	j	i	1	-i	-1

	C_0	\bar{C}_0	C_x	C_y	C_z
C_0	C_0	\bar{C}_0	C_x	C_y	C_z
\bar{C}_0	\bar{C}_0	C_0	C_x	C_y	C_z
C_x	C_x	C_x	$2C_0 + 2\bar{C}_0$	$2C_z$	$2C_y$
C_y	C_y	C_y	$2C_z$	$2C_0 + 2\bar{C}_0$	$2C_x$
C_z	C_z	C_z	$2C_y$	$2C_x$	$2C_0 + 2\bar{C}_0$

The 2π Point Defect is not the Trivial Defect

- One can imagine two line defects (generated by C_i for example) forming a 2π defect or annihilating each other
- The question: When will two π defects merge to create a 2π defect versus annihilating to a trivial defect?



Ex. Two z-Disclinations

- Case 1: no other disclinations
 - Combines to a 2π point defect

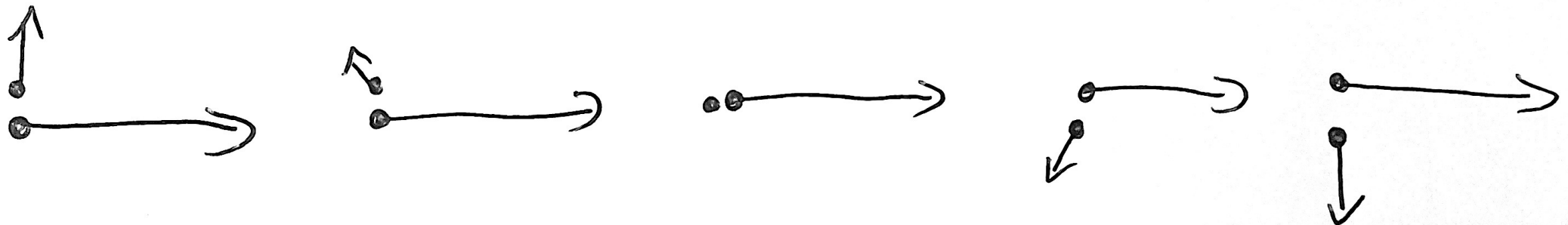
$$-1 = (i\sigma_z)(i\sigma_z)$$

- Case 2: one other disclination
 - Decays via “catalysis”

$$-(i\sigma_z) = (i\sigma_x)(i\sigma_z)(i\sigma_x)^{-1}$$

$$1 = (i\sigma_z)(-i\sigma_z)$$

Schematic of part of the decay process



Overall Guidance

- Based \rightarrow fundamental group elements

- The order matters

$$(i\sigma_x)(i\sigma_z)(i\sigma_z) \neq (i\sigma_z)(i\sigma_x)(i\sigma_z)$$

- Un-based \rightarrow conjugacy classes

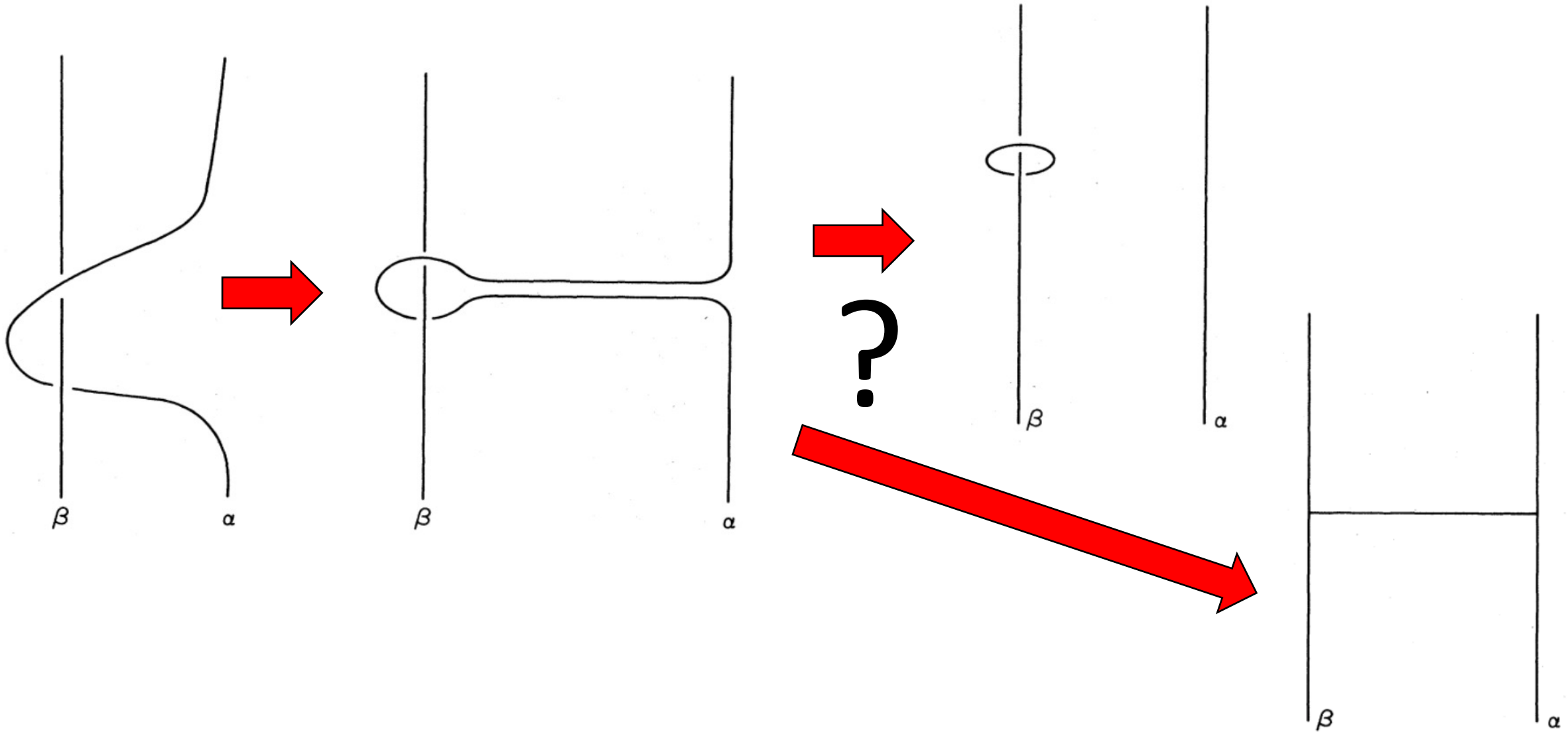
- Conjugacy classes can be ambiguous

- This ambiguity is lifted by the contour and what is within it
- Conjugacy classes combine in an Abelian fashion
- Think about the example from the last slide

Scarring: the Unique Feature of NA Groups

- For Abelian groups the order line defects cross doesn't matter
- For non-Abelian groups “scarring” can occur
 - An effect where the crossing of line defects generates another line defect
- Scarring is only avoided when $\beta\alpha\beta^{-1}\alpha^{-1} = 1$
 - But this only holds for commuting elements α and β
- Let's look at this pictorially

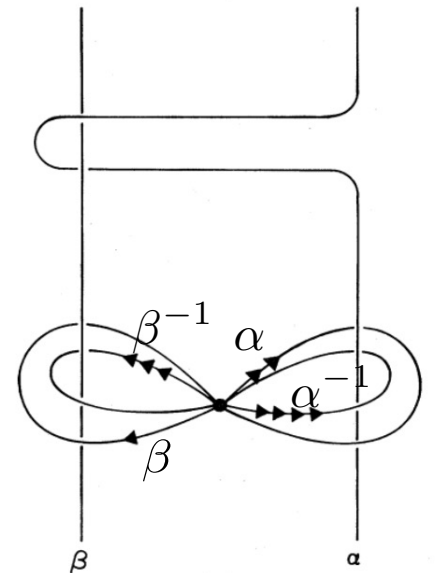
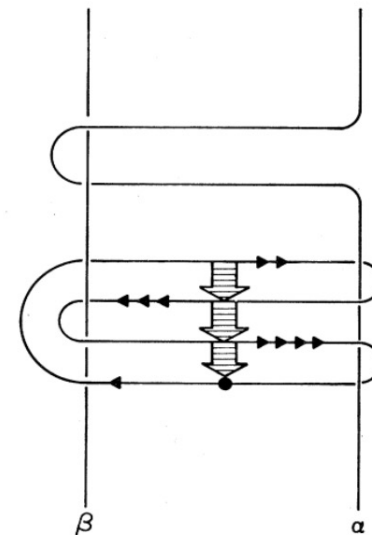
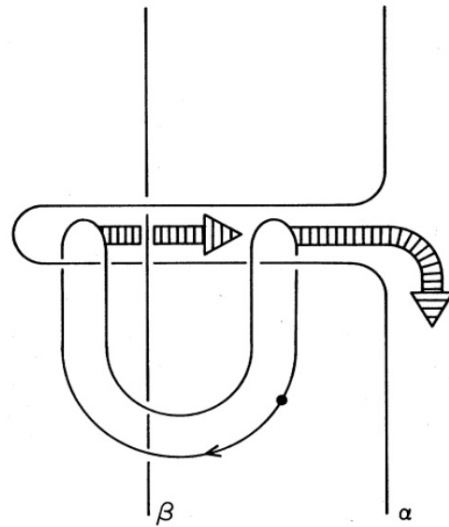
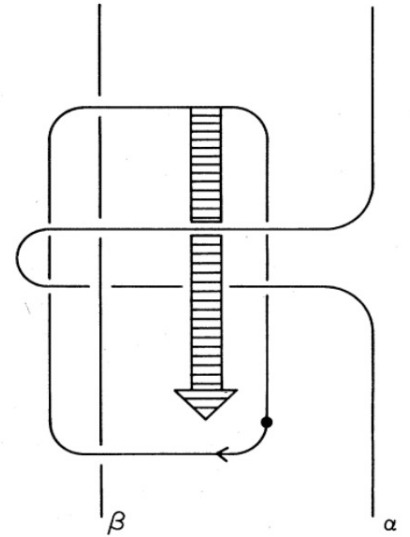
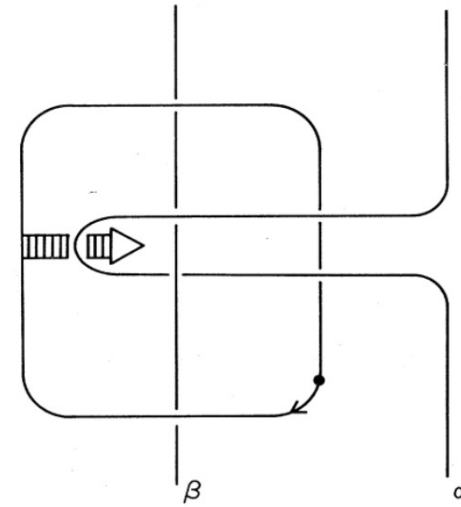
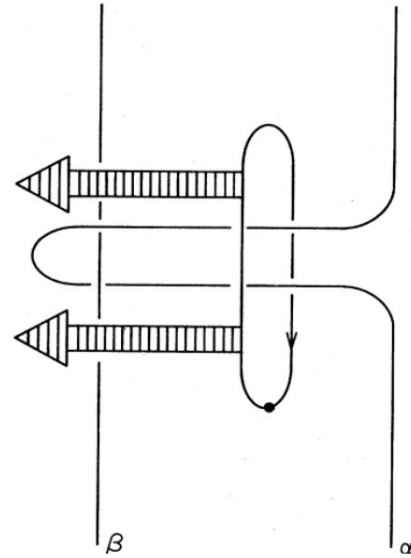
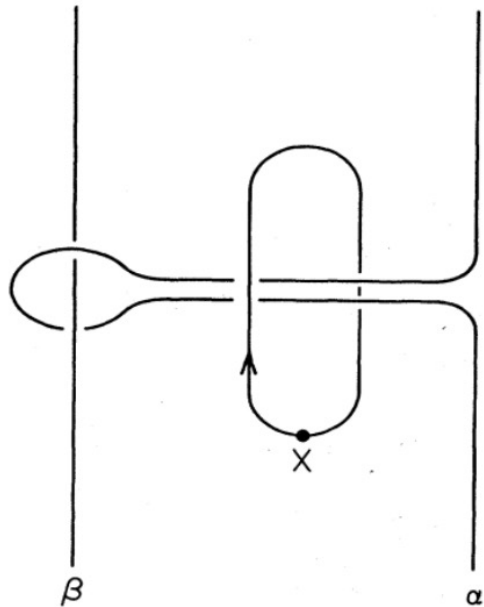
To Scar or not to Scar, that is the Question



Pictorial Approach to Smooth Deformations

- If we take our contour around the crossing we can obtain the condition for scarring

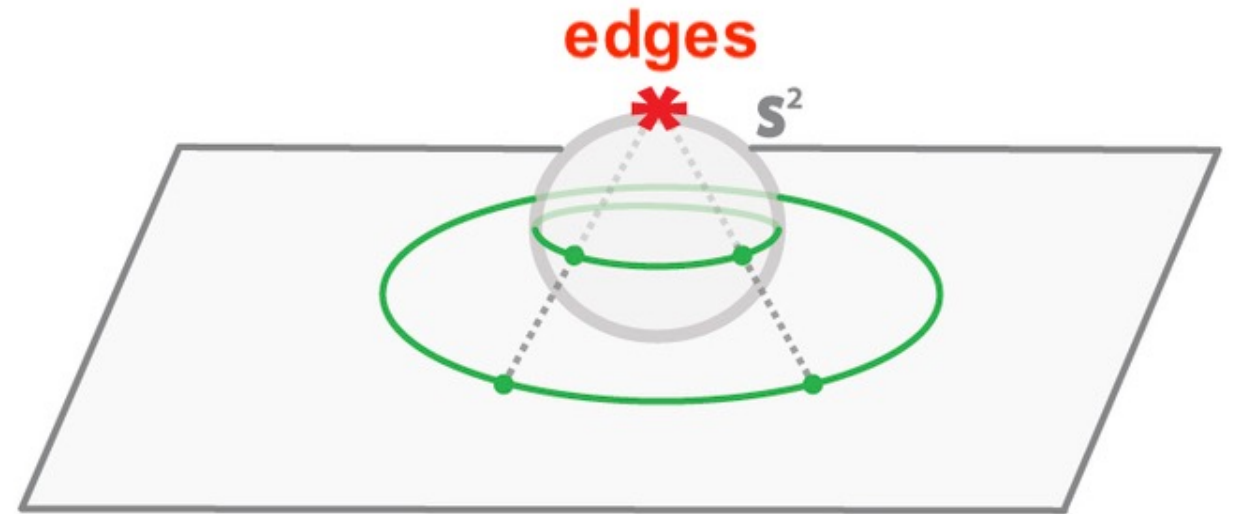
$$\beta\alpha\beta^{-1}\alpha^{-1} \neq 1$$



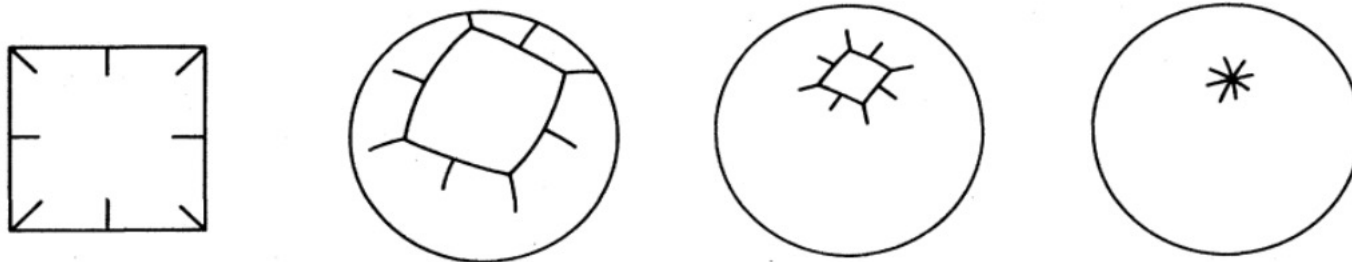
VII. The Second Homotopy Group

Spheres from Squares: 1 pt Compactification

- The one point compactification of the square is the sphere
- This will enable a useful notation where spheres with the same base point are connected at the edges
- “Closing the purse”

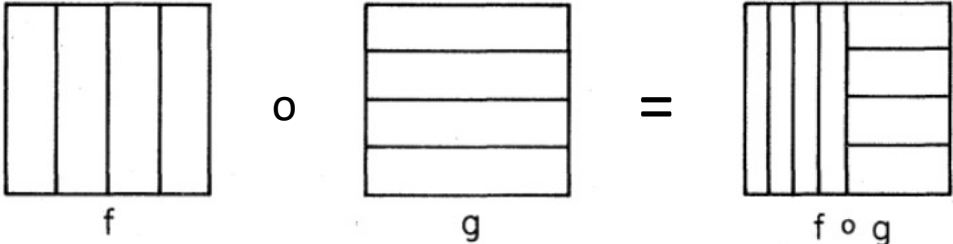


[Image Credit: Adapted from Antoniou and Lambropoulou](#)

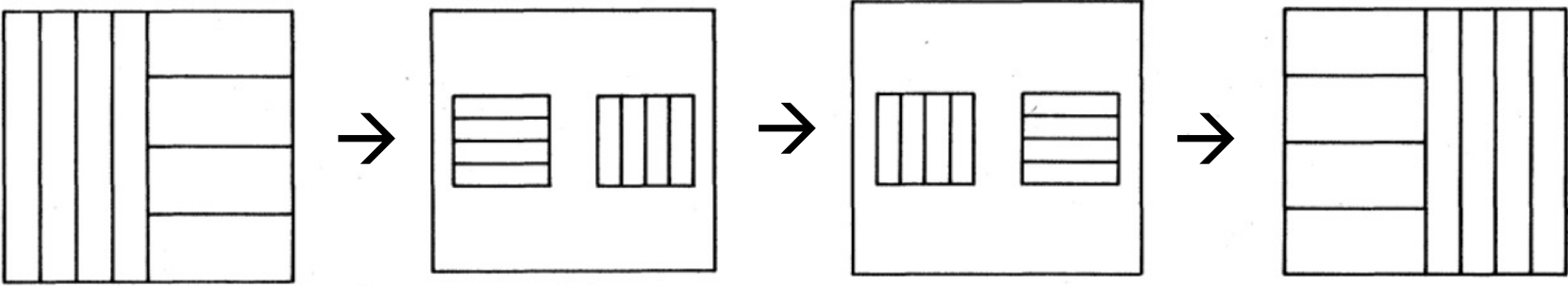


Theorem: π_2 is Abelian

- Composition of functions

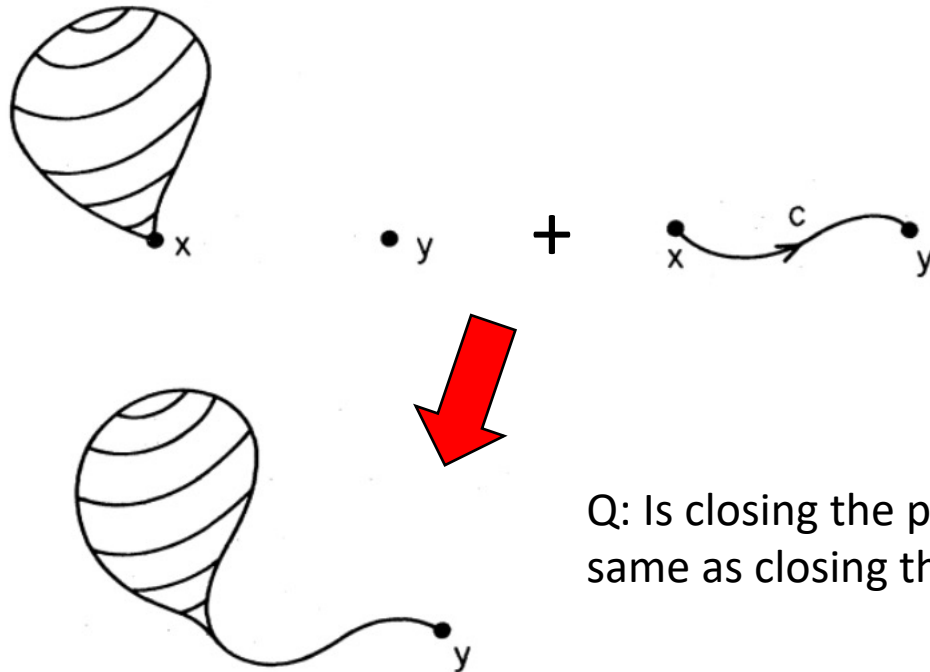


- Proof:

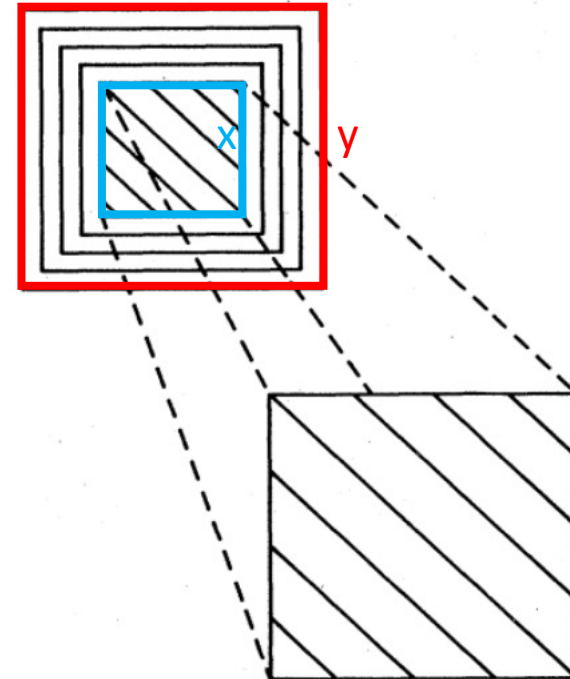


Theorem: π_2 is Base Point Independent

- Provided the spheres include the same volume
- Shifting to other base points requires that the enclosed region be **2-simple**, i.e. c induces a trivial automorphism on π_2



Q: Is closing the purse here the same as closing the purse there?



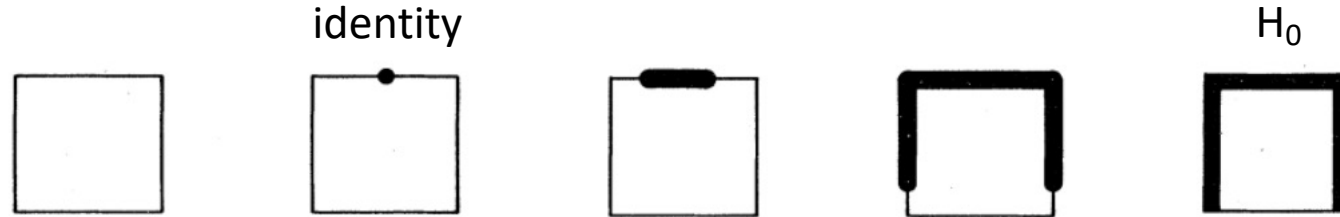
Fundamental Theorem of π_2

- Let G be simply connected (always possible—make it big enough)
- Let $\pi_2(G)=0$ (ok by Cartan's theorem for compact Lie groups)
- Let H be the isotropy subgroup ($H = \{g \text{ in } G \text{ such that } g \Psi = \Psi\}$)
- The coset space is G/H
- Let H_0 be the connected component connected to the identity
- Then **$\pi_2(\mathbf{G/H}) = \pi_1(\mathbf{H_0})$**

- Proof: using the assumptions above establish an isomorphism between the two—one to one, and same algebraic structure

Pictorial Interpretation of the Fundamental Thm.

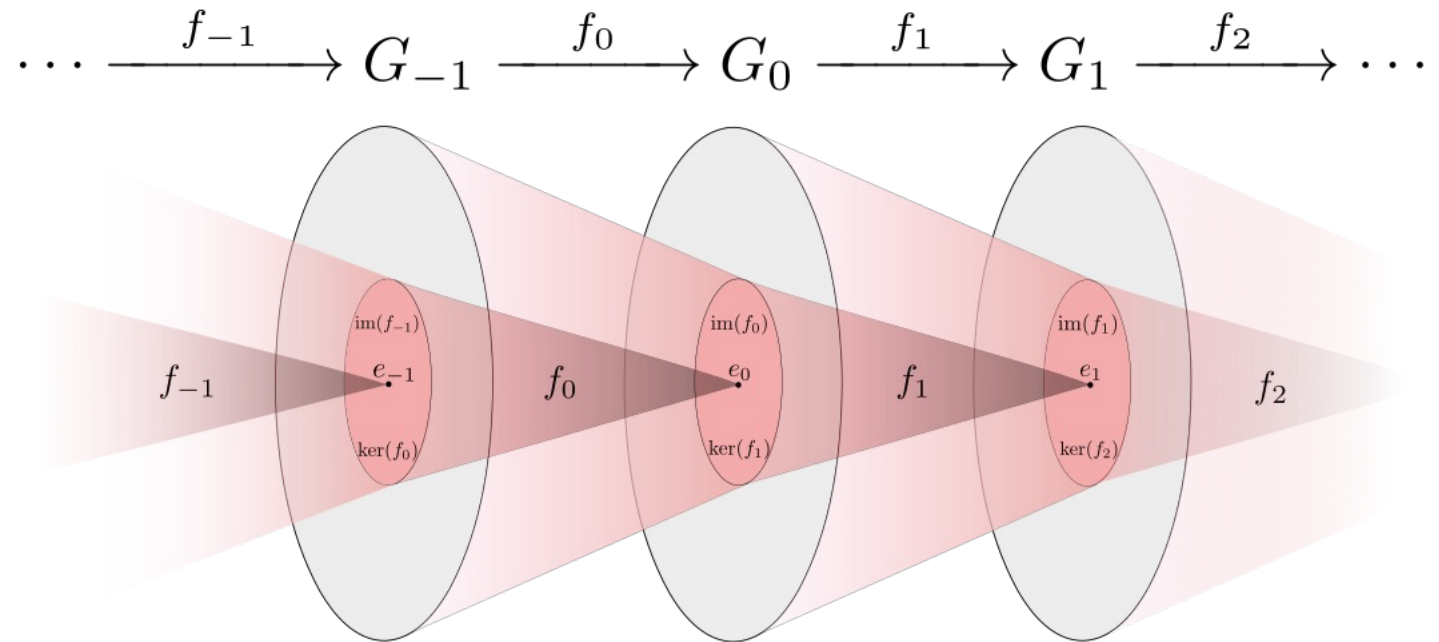
- Structure of H_0 on the edge of π_2 of the region



- Windings about H_0 gives us the possible “mouths” of the “purses”
- Here the $\pi_2(G)$ is assumed to be trivial so the only non-trivial behavior comes from H_0 and its winding about the “purse”

Exact Sequences

- This is an example of an exact sequence
- Since $\pi_2(G) = 0$ we can map to the kernel of H (H_0)
- This simplifies the computation of higher homotopy groups
- Exact sequences lead to the Bott periodicity theorem



[Image credit: Wikipedia](#)

Corollary

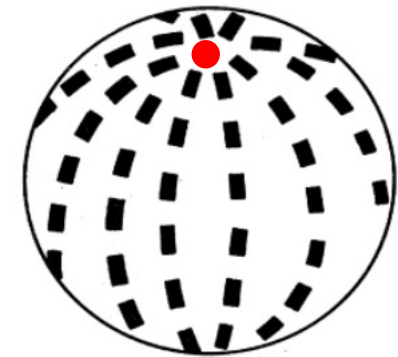
- If H is discrete, then H_0 is a simple point
- Hence $\pi_1(H_0) = 0$
- So that $\pi_2(G/H) = \pi_1(H_0) = 0$

- Consequence

- H is discrete for biaxial nematics and superfluid ^3He
- These systems do not exhibit point defects in 3D



NO BIAXIAL HEDGEHOGS!



Ex. Spins and Nematics

- Spins

- $G = SU(2)$
- $H = \{\text{rotations about one axis}\}$
- $H_0 = U(1)$
- $\pi_2(G/H) = \pi_1(H_0) = \pi_1(U(1)) = \{\text{winding numbers}\} = \mathbb{Z}$

- Nematics

- The same but an additional symmetry
 - $+n$ and $-n$ winding numbers correspond to the same field configuration
- $\pi_2(G/H) = \mathbb{Z}^+$

Comparison of π_1 and π_2

System	π_1 (First Homotopy Group)	π_2 (Second Homotopy Group)
Planar spins	\mathbb{Z}	Not sensible
Spins in 3D	0	\mathbb{Z}
Nematics	$\mathbb{Z}/2\mathbb{Z}$	\mathbb{Z}^+
Biaxial nematics	Quaternions	0
Superfluid ^3He	$\mathbb{Z}/2\mathbb{Z}$	0

Bonus. Applications to Q. Condensed Matter

Ex. Chern Numbers

- Consider a $N \times N$ Hamiltonian h (Hermitian matrix)
- We can decompose as $h = U D U^{-1}$
- Computation
 - $G = U(N)$
 - $H = DU(N)$
 - $H_0 = SDU(N)$
 - $\pi_2(G/H) = \pi_1(H_0) = \pi_1(SDU(N)) = \mathbb{Z}^N/\mathbb{Z} = \mathbb{Z}^{N-1}$
- Interpretation: N bands with integer “Chern” numbers c_n and $\sum c_n = 0$
- (for more details and related discussion see [Moore, page 9](#))

Bott Periodicity Theorem

- From structure of exact sequences

- Unitary Groups

- For $N \geq (n+1)/2$

$$\pi_n(U(N)) \cong \pi_n(SU(N)) \cong \begin{cases} 1 & n \text{ even} \\ \mathbb{Z} & n \text{ odd} \end{cases}$$

- Orthogonal Groups

- For $N \geq n+2$

$$\pi_n(O(N)) \cong \pi_n(SO(N)) \cong \begin{cases} 1 & n \equiv 2, 4, 5, 6 \pmod{8} \\ \mathbb{Z}/2\mathbb{Z} & n \equiv 0, 1 \pmod{8} \\ \mathbb{Z} & n \equiv 3, 7 \pmod{8} \end{cases}$$

Periodic Table of Topological Insulators (& SC)

- This Bott periodicity appears shows up in the classification of (quadratic) mean field topological insulators/superconductors

Class	T	C	S	δ							
				0	1	2	3	4	5	6	7
A	0	0	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0
AIII	0	0	1	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}	0	\mathbb{Z}
AI	+	0	0	\mathbb{Z}	0	0	0	$2\mathbb{Z}$	0	\mathbb{Z}_2	\mathbb{Z}_2
BDI	+	+	1	\mathbb{Z}_2	\mathbb{Z}	0	0	0	$2\mathbb{Z}$	0	\mathbb{Z}_2
D	0	+	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0	$2\mathbb{Z}$	0
DIII	-	+	1	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0	$2\mathbb{Z}$
AII	-	0	0	$2\mathbb{Z}$	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0	0
CII	-	-	1	0	$2\mathbb{Z}$	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0	0
C	0	-	0	0	0	$2\mathbb{Z}$	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}	0
CI	+	-	1	0	0	0	$2\mathbb{Z}$	0	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}

Image credit: Pulcu (2021)

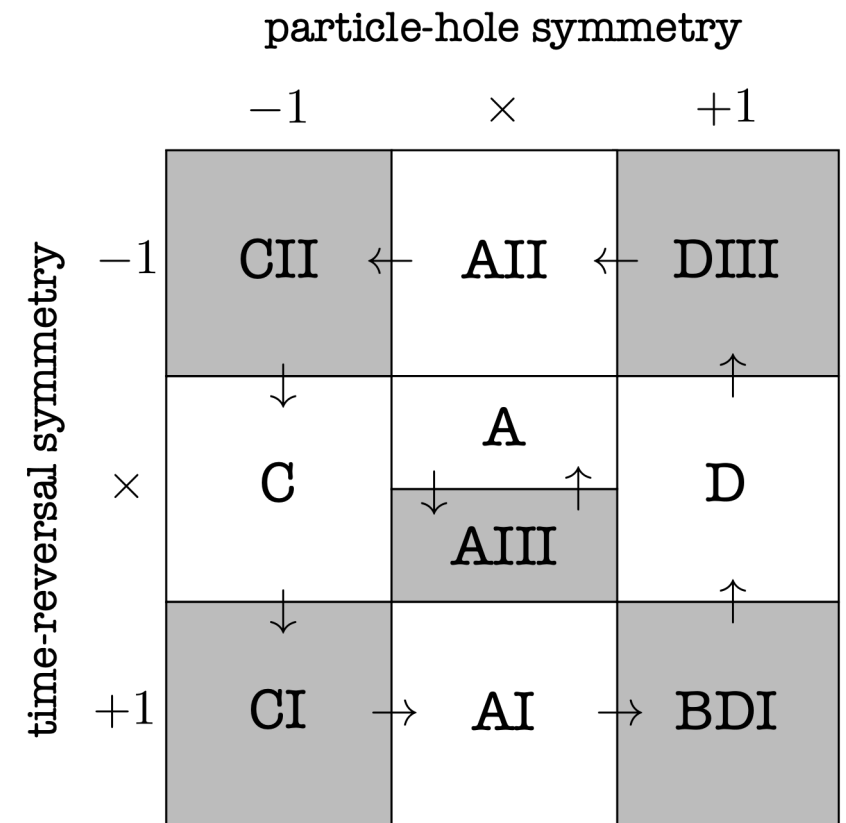


Image credit: Topology in Condensed Matter

Some Homotopy Groups

Table 4.1. Useful homotopy groups.

	π_1	π_2	π_3	π_4	π_5	π_6
SO(3)	\mathbb{Z}_2	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}
SO(4)	\mathbb{Z}_2	0	$\mathbb{Z} + \mathbb{Z}$	$\mathbb{Z}_2 + \mathbb{Z}_2$	$\mathbb{Z}_2 + \mathbb{Z}_2$	$\mathbb{Z}_{12} + \mathbb{Z}_{12}$
SO(5)	\mathbb{Z}_2	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	0
SO(6)	\mathbb{Z}_2	0	\mathbb{Z}	0	\mathbb{Z}	0
SO(n) $n > 6$	\mathbb{Z}_2	0	\mathbb{Z}	0	0	0
U(1)	\mathbb{Z}	0	0	0	0	0
SU(2)	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}
SU(3)	0	0	\mathbb{Z}	0	\mathbb{Z}	\mathbb{Z}_6
SU(n) $n > 3$	0	0	\mathbb{Z}	0	\mathbb{Z}	0
S^2	0	\mathbb{Z}	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}
S^3	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2	\mathbb{Z}_{12}
S^4	0	0	0	\mathbb{Z}	\mathbb{Z}_2	\mathbb{Z}_2
G_2	0	0	\mathbb{Z}	0	0	\mathbb{Z}_3
F_4	0	0	\mathbb{Z}	0	0	0
E_6	0	0	\mathbb{Z}	0	0	0
E_7	0	0	\mathbb{Z}	0	0	0
E_8	0	0	\mathbb{Z}	0	0	0

From Nakahara (2003)

References

- N.D. Mermin, *Topological theory of defects*, RMP 51, 591 (1979)
- J.E. Moore, *Notes for MIT minicourse on topological phases* (2011)
- Nakahara, *Geometry, Topology, and Physics*, 2nd Edition (2003)
- [Topology in Condensed Matter \(2015\)](#)