

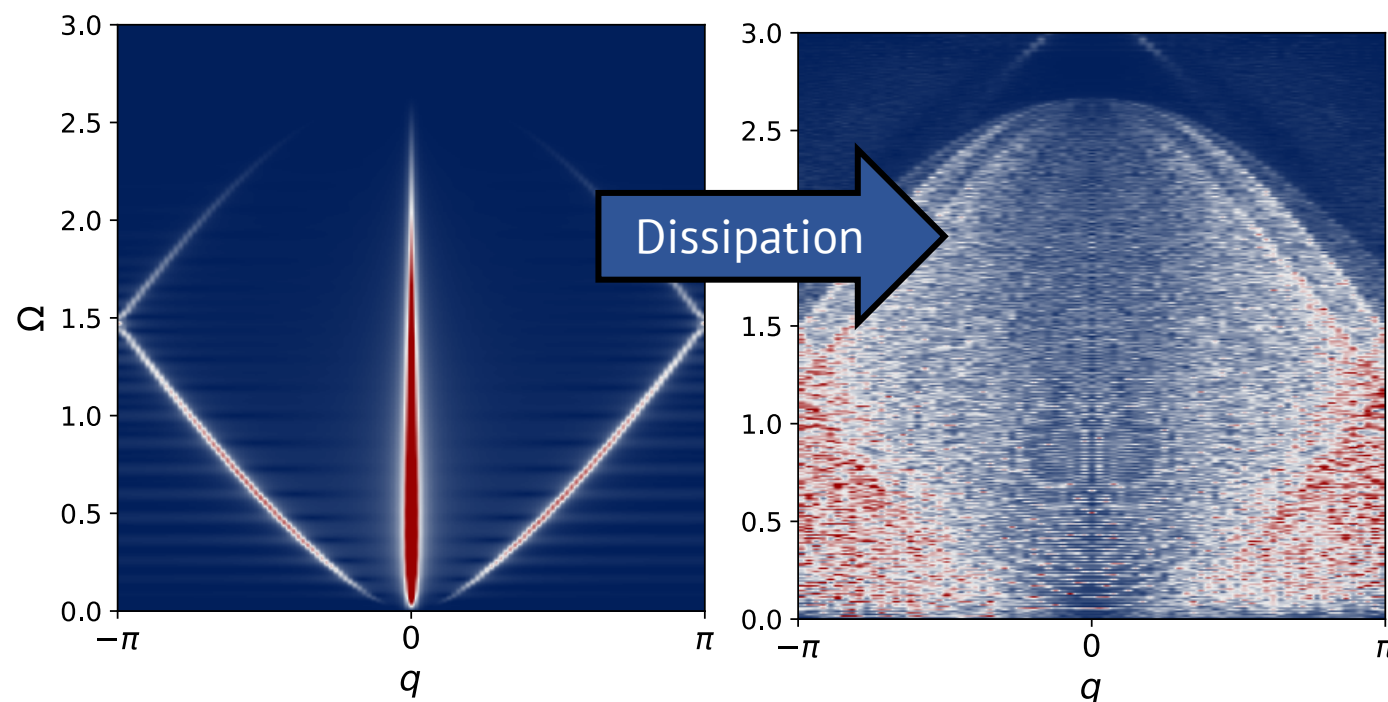
Dynamic Response of Dissipative Spin Chains

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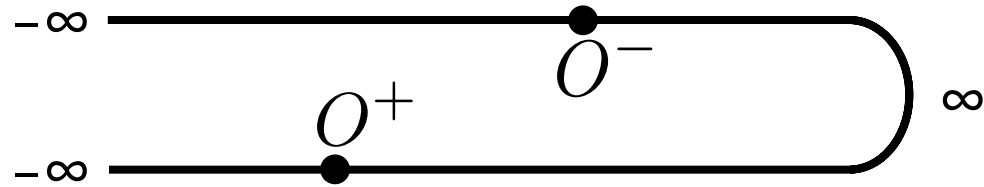
University of Pennsylvania

5 March 2024

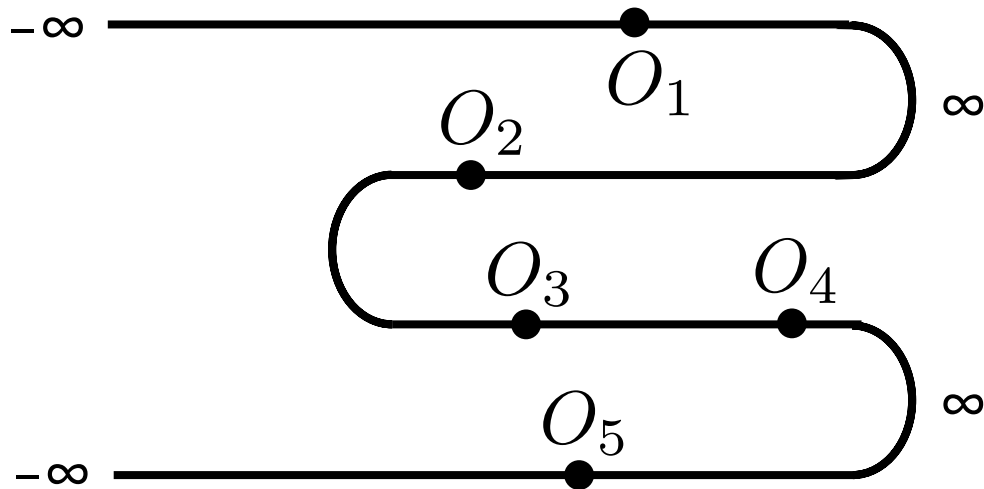


Contours and Correlators

- Take expectation values at times on the Keldysh contour



- Can express two-point correlators in terms of G^R , G^A and G^K
- Multiple contours is generalized by OTOCs



Lindblad Master Equation

- Time evolution in the limit of continuous measurement by a memoryless bath $i\dot{\rho} = \mathcal{L}[\rho]$

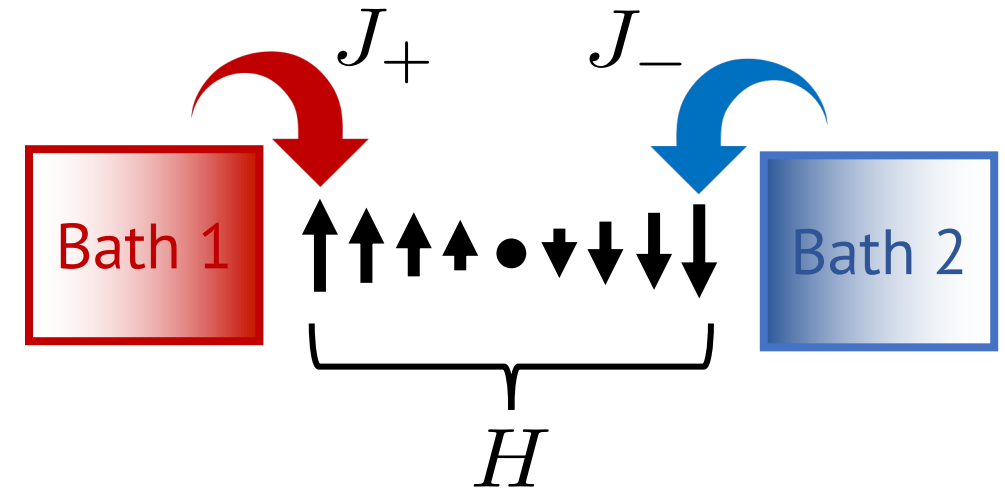
$$i\dot{\rho} = (\mathcal{L}^{\text{coh.}} + \mathcal{L}^{\text{n.h.}} + \mathcal{L}^{\text{jump}})[\rho]$$

where,

$$\mathcal{L}^{\text{coh.}}[\rho] = [H, \rho]$$

$$\mathcal{L}^{\text{n.h.}}[\rho] = -i\frac{\Gamma}{2} \sum_m \{J_m^\dagger J_m, \rho\}$$

$$\mathcal{L}^{\text{jump}}[\rho] = i\frac{\Gamma}{2} \sum_m 2J_m \rho J_m^\dagger$$



Spin chain with boundary dissipation. Jump operators linear in spins map to jump operators linear in fermions under Jordan-Wigner transformation.

Energy scales

J: spin-spin coupling

h: transverse field

Γ : dissipation strength

Green's Functions

- We want to probe the steady state with static density response $\langle S^z \rangle_{ss}$ and dynamic response $\mathcal{S}_{i,j}^z(\Omega) = \int dt e^{-i\Omega t} \langle [S_i^z(t), S_j^z(0)] \rangle_{ss} \theta(t)$

$$\langle S^z \rangle_{ss} = \Omega \text{ [loop diagram with } G^K \text{ and } S^z \text{]} \Omega$$

$$\mathcal{S}_{i,j}^z(\Omega) = \Omega \text{ [loop diagram with } G^R, G^K, S_j^z, S_i^z \text{]} \Omega + \Omega \text{ [loop diagram with } G^A, G^K, S_j^z, S_i^z \text{]} \Omega$$

- These can be expressed in terms of Green's functions
 - Closed system: G^R and G^A
 - Open system: G^R , G^A and a Keldysh Green's function G^K
- Jump terms lead to complex self energy!
- Wick's theorem can be used to reduce multi-point correlation functions to two-point correlation functions

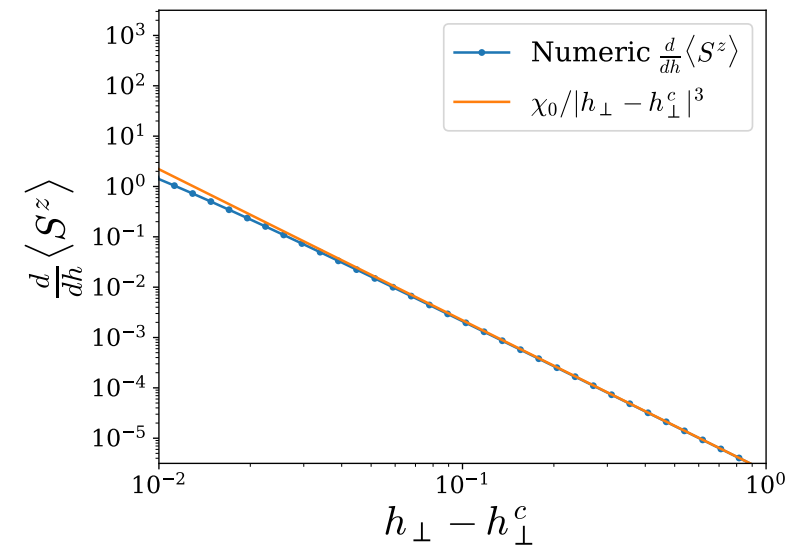
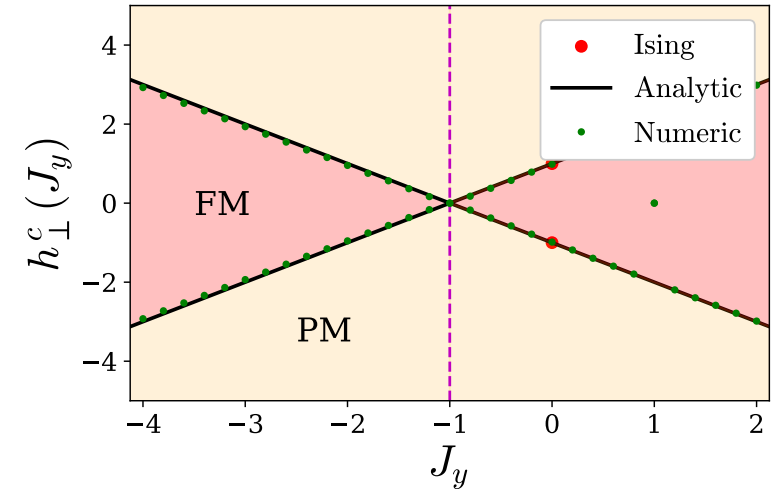
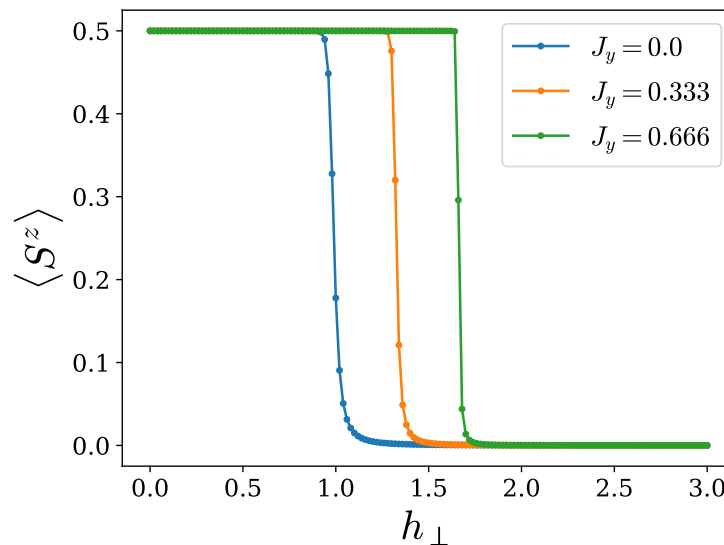
Lindblad-Keldysh GFs:
 • Thompson and Kamenev, Ann. Phys. 455, 169385 (2023)
 • McDonald and Clerk, Phys. Rev. Res. 5, 033107 (2023)

Ising and XY Models: Closed Systems I

- Transverse field XY spin chain is paradigmatic example of gapped/gapless quantum matter

$$H_{XY} = \sum_{n=1}^{N-1} J_n^x S_n^x S_{n+1}^x + J_n^y S_n^y S_{n+1}^y + \sum_{n=1}^N h_n S_n^z$$

- Jordan-Wigner transformation to fermions
- Find magnetization
 - Paramagnetic to ferromagnetic transition
- Susceptibility diverges with a power law

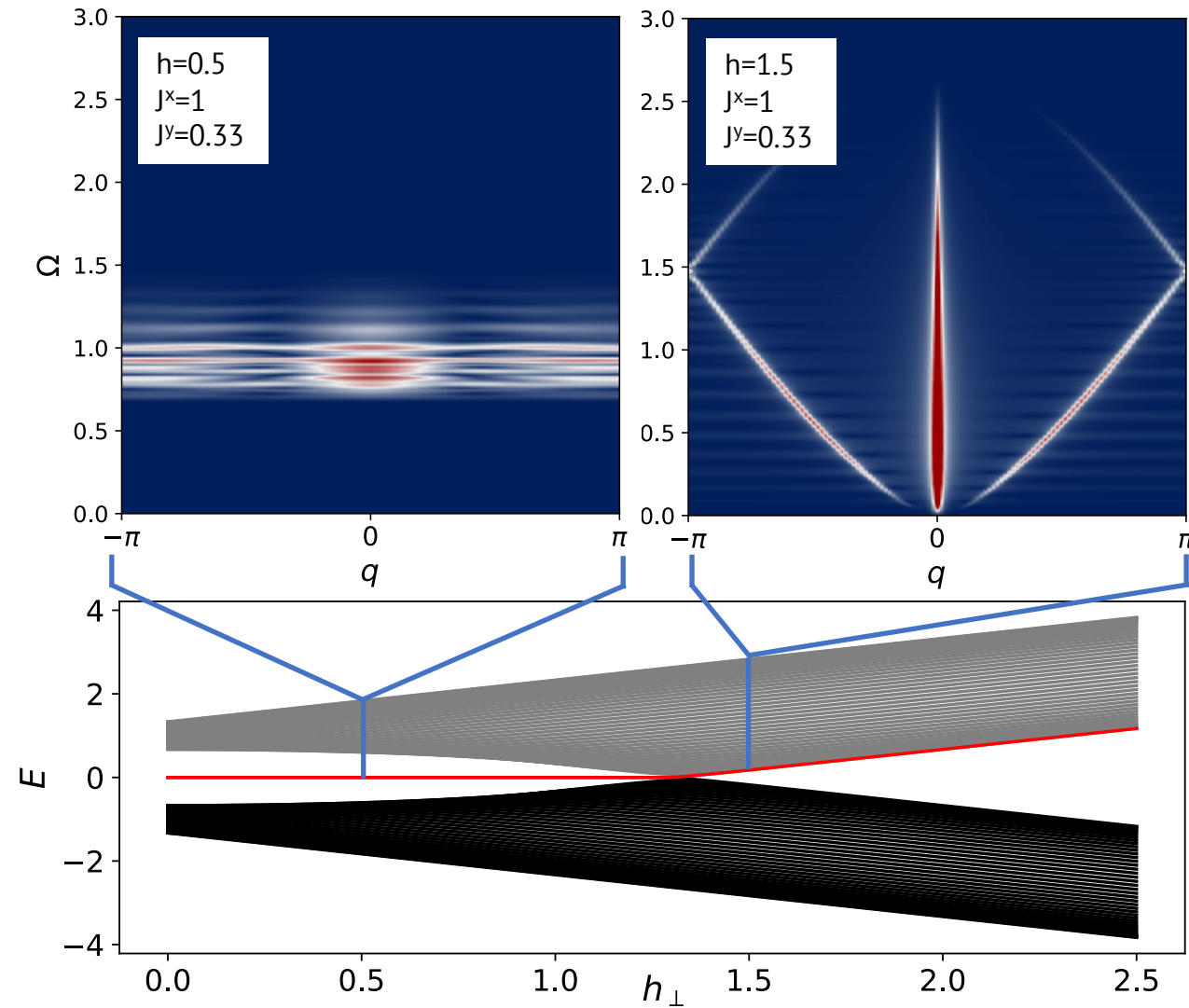


Ising and XY Models: Closed Systems II

- Dynamic correlation function

$$\chi_{ij}(\Omega) = \sum_{n>0} \frac{\langle u_0 | S_i^z | u_n \rangle \langle u_n | S_j^z | u_0 \rangle}{\Omega + i\eta - (\epsilon_n - \epsilon_0)}$$

- Probe excitations above $|u_0\rangle$
- See the structure of gapped and gapless excitations
- All negative energy states are filled (think Pauli blocking)



Closed Systems

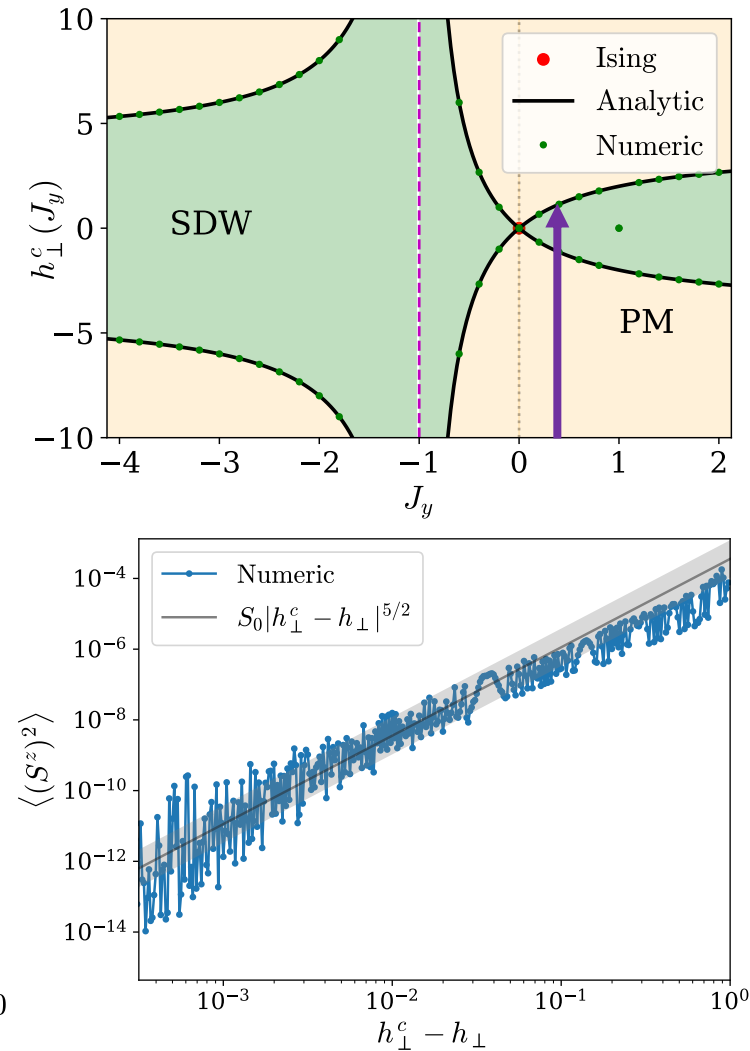
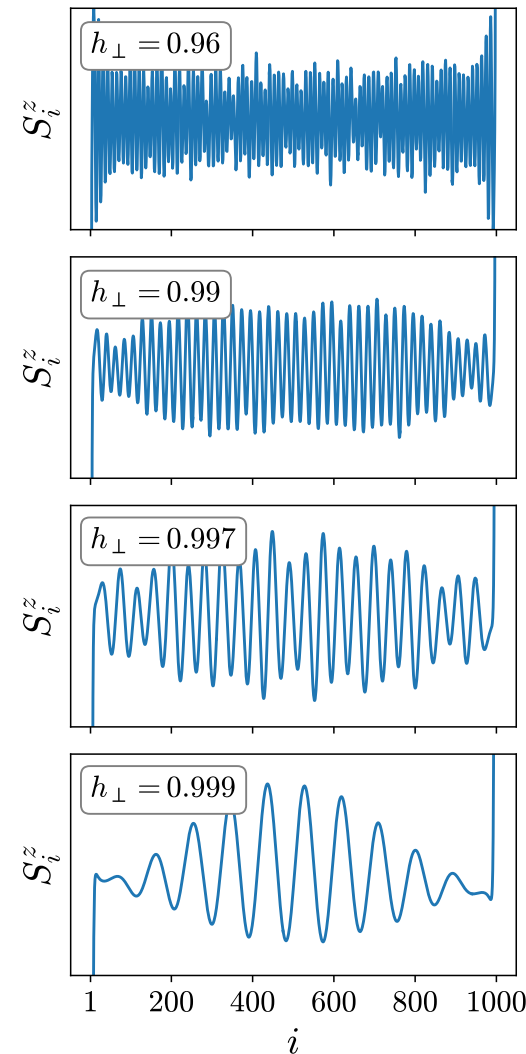
- Hermitian
- Real eigenvalues
- Orthogonal eigenstates
- Fermi-Dirac and Bose-Einstein distributions

Open Systems

- Non-Hermitian
- Complex eigenvalues
- Bi-orthogonal eigenstates
- Distribution function given by G^K
 - In Lindblad formalism this does not recover FD and BE distribution functions in the limit of weak dissipation

Ising and XY Models: Open Systems I

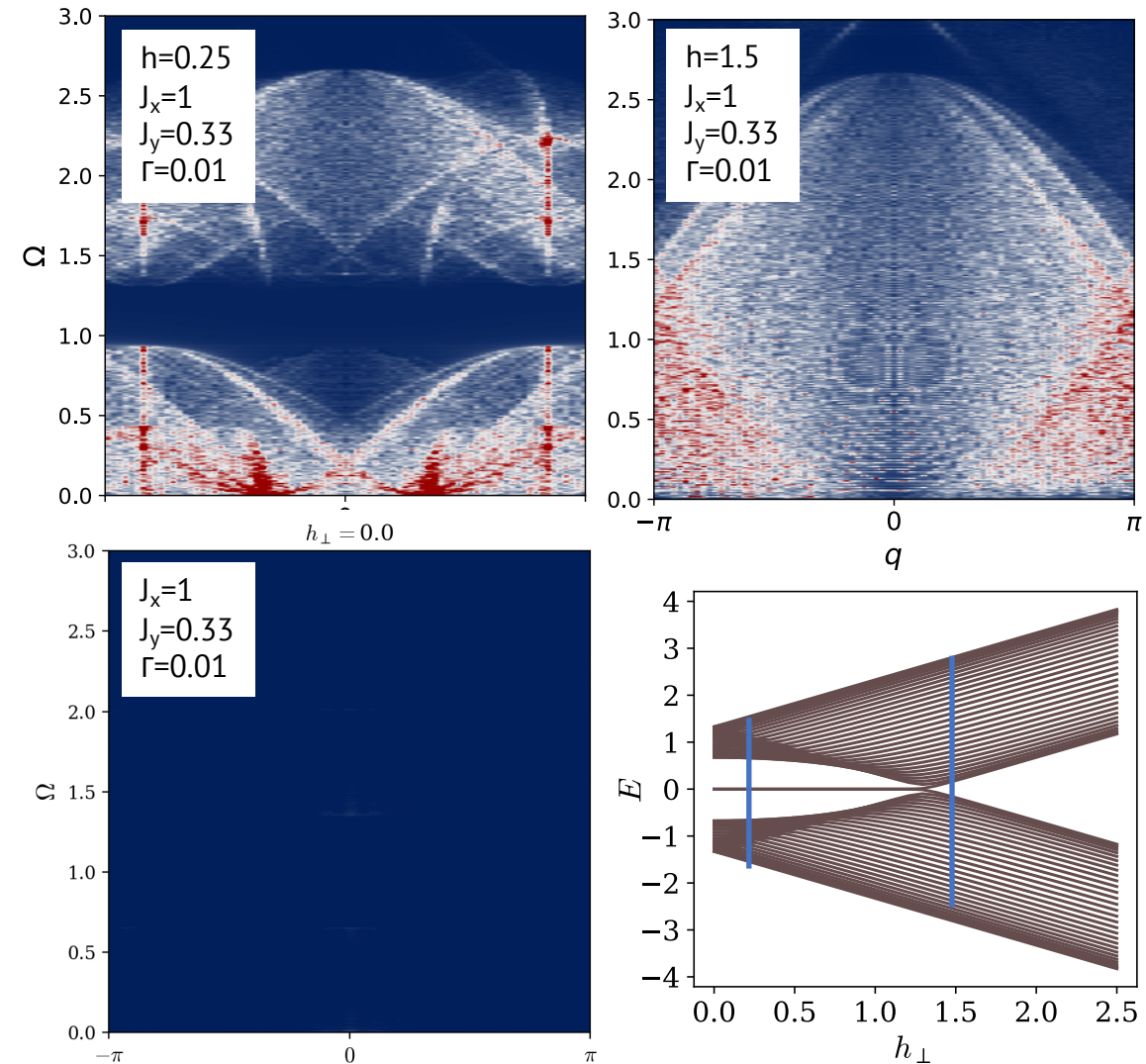
- Spin chain with boundary dissipation S^+ on left, S^- on right
- New phase: spin-density wave
 - Magnetization $\langle S^z \rangle$ vanishes, but its higher moments do not
 - Phase boundary is very different from equilibrium phase boundary
- Wavelength diverges and $\langle (S^z)^2 \rangle$ exhibits power-law scaling on near critical point h^c
- Choose model with $J^x=1, J^y=1/3$ so that $h^c = 1$



- Dynamic correlation function

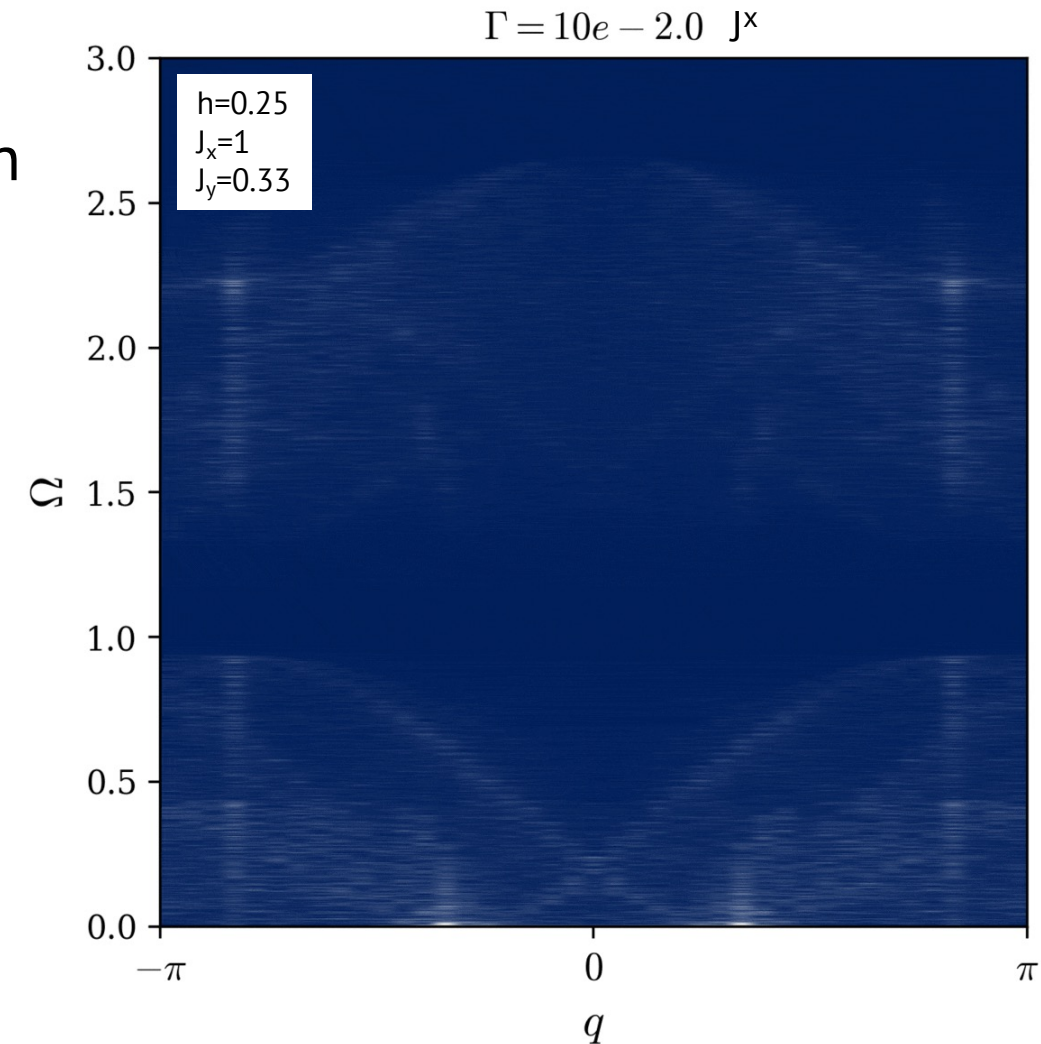
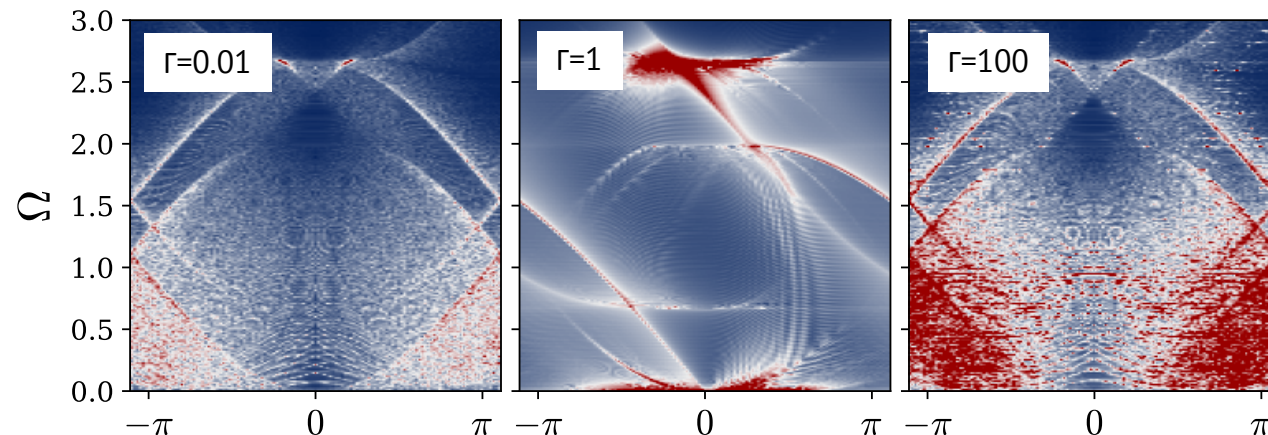
$$S_{i,j}^z(\Omega) = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \left(\text{Tr}[S_i^z G^R(\omega) S_j^z G^K(\omega + \Omega)] + \text{Tr}[S_i^z G^K(\omega - \Omega) S_j^z G^A(\omega)] \right)$$

- Fractional occupation given by distribution function for ρ_{ss}
 - States at all energies contribute
 - Spectral gaplessness for $h < h^c$
- Vertical lines at q_{SDW} from SDW
- Dispersing modes from $2q_{SDW}$



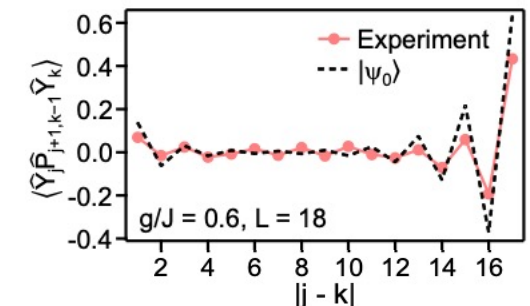
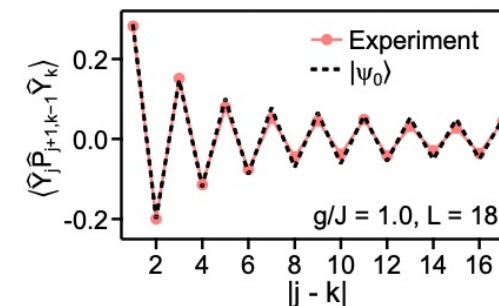
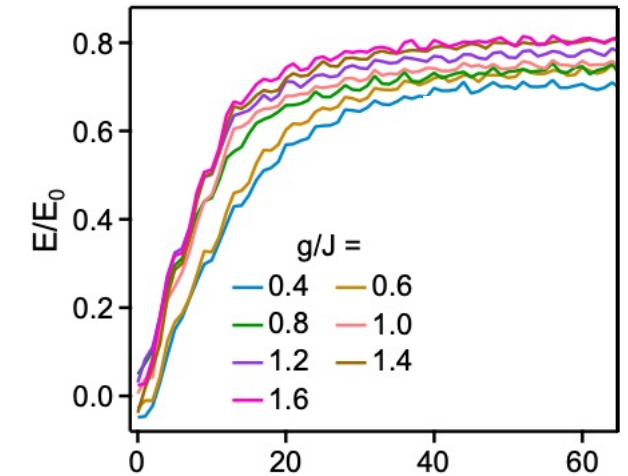
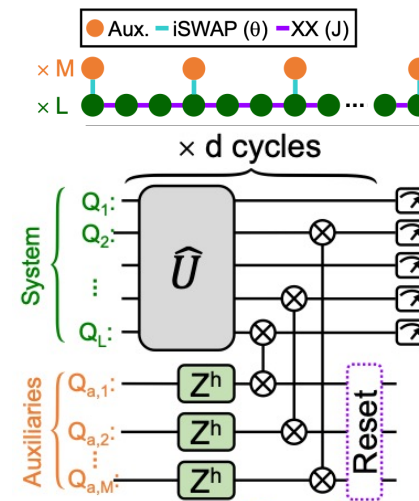
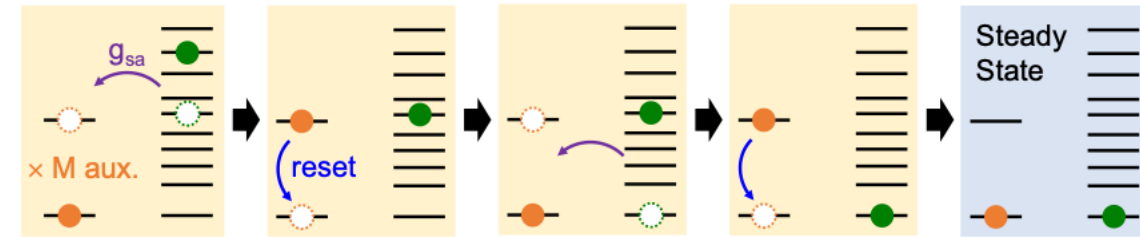
Ising and XY Models: Open Systems III

- Strong coupling ($\Gamma \sim J$)
 - Inversion-symmetry broken by dissipation
 - Large amplitude response
- Weak ($\Gamma \ll J$) and ultra-strong ($\Gamma \gg J$)
 - Restoration of inversion symmetry
 - Bulk is relatively isolated from boundaries
 - Small amplitude response



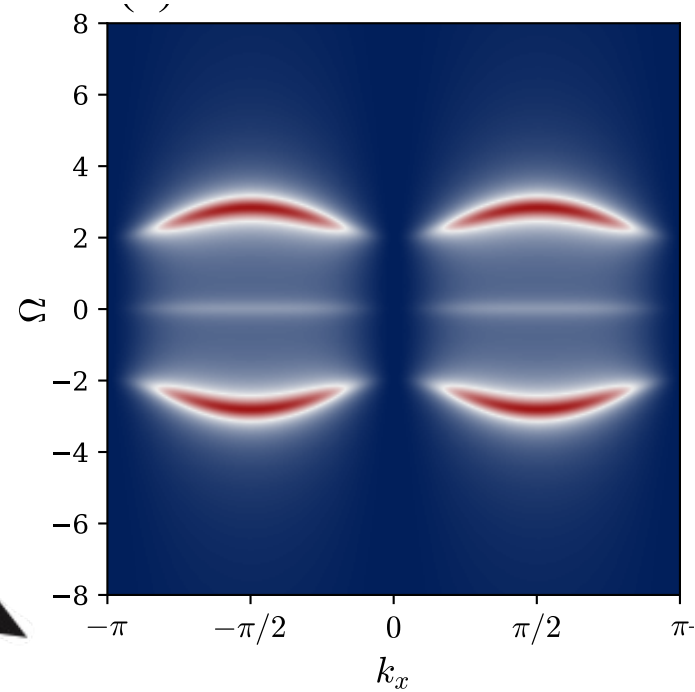
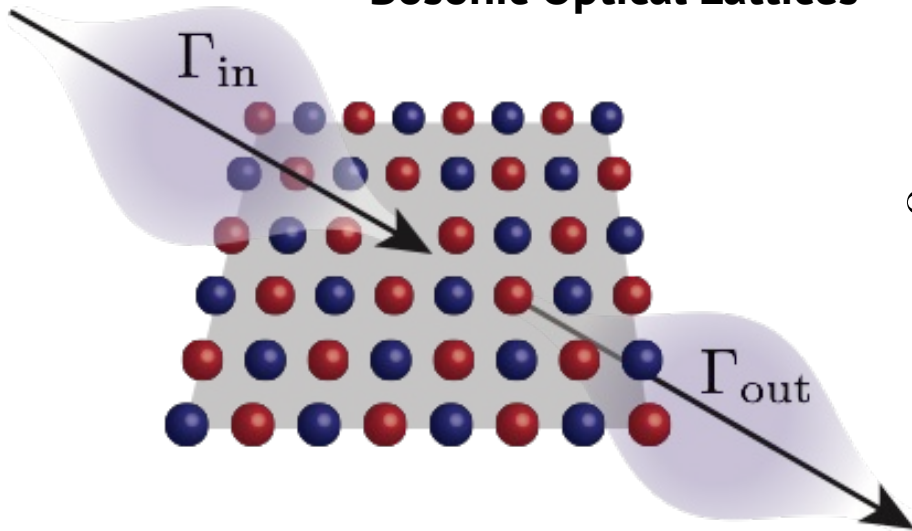
Dissipation on NISQ Hardware

- Google QuantumAI experiment
 - Dissipative Ising model (2D)
 - Use ancillas with resets to simulate dissipative memoryless baths
 - Mi *et al*, 2304.13878
- Cool the system towards its ground state through stroboscopic evolution: unitary evolution followed by dissipative evolution
- Potential to realize other strongly dissipative systems on NISQ devices

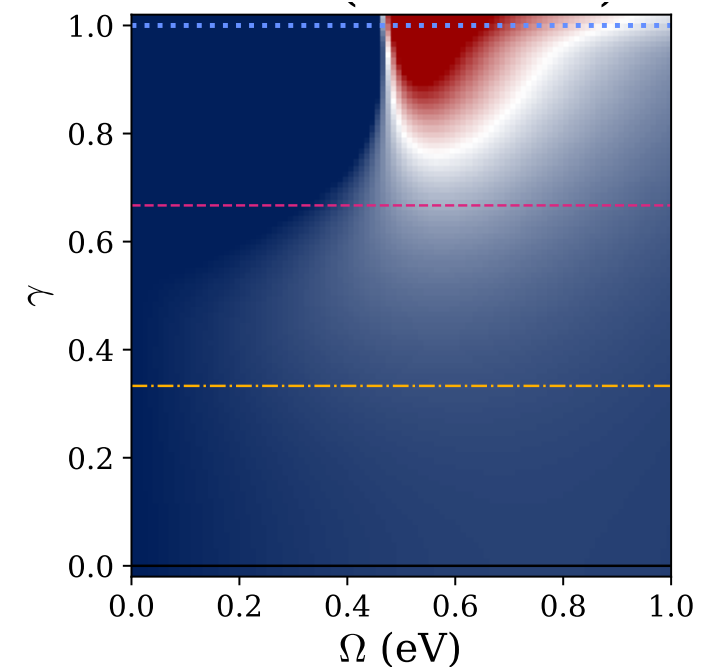


- Accessible introduction to Lindblad-Keldysh Green functions
- Expressions to directly calculate responses

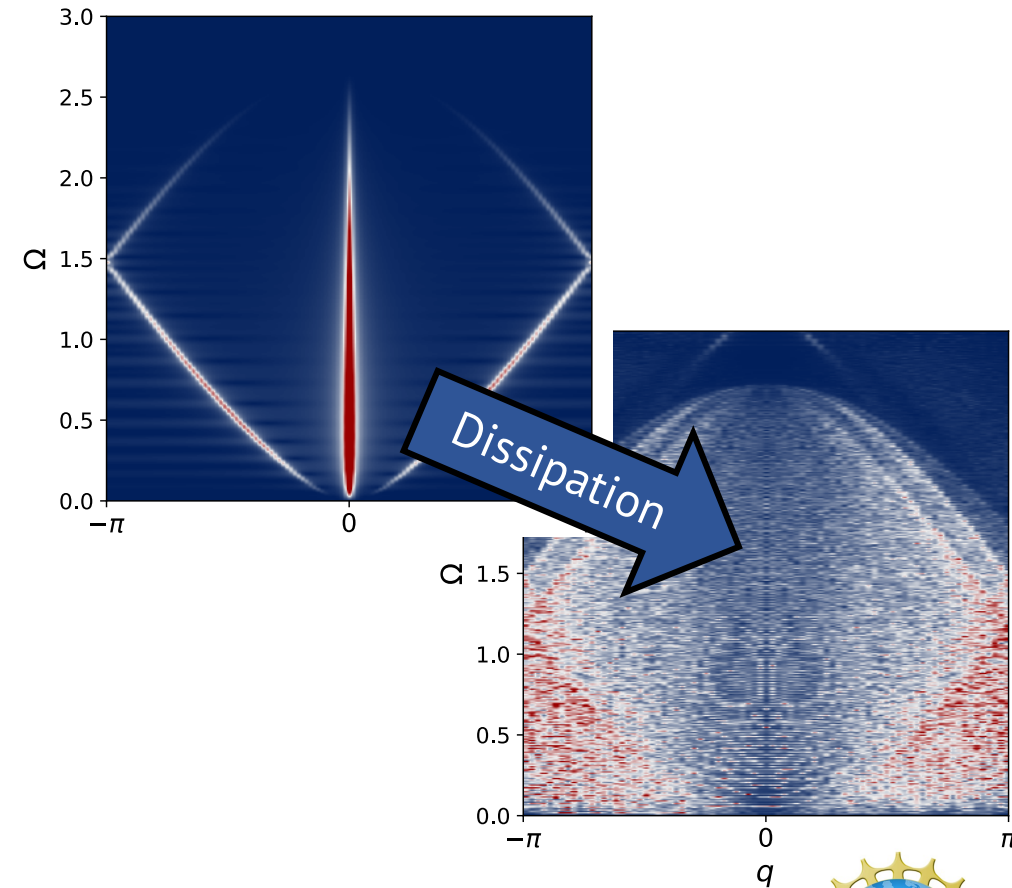
Dynamic Response in
Bosonic Optical Lattices



Nonlinear Shift Response in
Bernal Bilayer Graphene



- Dissipative engineering is a promising direction to control systems and realize out of equilibrium physics
- Finite frequency probes of excitations are readily accessible in our formalism
- Generalizations to interacting systems are the frontier
 - Prethermal plateaus and avoiding thermalization at long-times
 - RPA and collective dissipative phenomena
 - And more!



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